Calibrationless rotating Lorentz-force flowmeters for low flow rate applications

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Calibrationless rotating Lorentz-force flowmeters for low flow rate applications

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Abstract

A ‘weighted magnetic bearing’ has been developed to improve the performance of rotating Lorentz-force flowmeters (RLFFs). Experiments have shown that the new bearing reduces frictional losses within a double-sided, disc-style RLFF to negligible levels. Operating such an RLFF under ‘frictionless’ conditions provides two major benefits. First, the steady-state velocity of the RLFF magnets matches the average velocity of the flowing liquid at low flow rates. This enables an RLFF to make accurate volumetric flow measurements without any calibration or prior knowledge of the fluid properties. Second, due to minimized frictional losses, an RLFF is able to measure low flow rates that cannot be detected when conventional, high-friction bearings are used. This paper provides a brief background on RLFFs, gives a detailed description of weighted magnetic bearings, and compares experimental RLFF data to measurements taken with a commercially available flowmeter.

Keywords: low-friction bearings, Lorentz-force velocimetry, flowmeter, liquid metal

1. Introduction

Shercliff developed the rotating Lorentz-force flowmeter (RLFF) in the 1950s [1–3]. As depicted in figure 1, a double-sided, disc-style RLFF consists of evenly-spaced permanent magnets that are installed near the periphery of a disc or flywheel. The center of the magnet assembly is connected to a low-friction bearing that permits rotational motion. When the flowmeter is installed alongside a pipe or tube that is filled with a flowing, electrically-conductive liquid, the resultant Lorentz-force between the liquid and magnets generates a torque upon the magnet assembly that causes it to rotate. During operation, the average velocity of the liquid can be determined by measuring the corresponding angular velocity of the flowmeter.

RLFFs are inexpensive to manufacture and simple to install. Moreover, RLFFs are non-contact devices that do not introduce any moving parts or seals into piping or tubing networks, so they can safely operate within systems containing chemically aggressive, hazardous, or very high-temperature fluids (e.g. molten metals, strong acids, etc). Collectively, these features make RLFFs useful instruments that could benefit the nuclear, concentrated solar power, chemical, pharmaceutical, and metallurgical/casting industries.

Previous work with RLFFs has shown that a minimum flowrate is required to overcome the static friction or ‘stiction’ found in conventional bearings [4–6]. Below the minimum flowrate, an RLFF does not spin and velocity measurements are impossible. Depending on the bearings used within an RLFF, the minimum flowrate could be quite large. Above the minimum flow rate, once the RLFF has been set in to motion, frictional forces acting upon the flowmeter must be accounted for in order to accurately correlate the angular velocity of the RLFF to the velocity of the flowing liquid [5, 7].

To avoid this issue, researchers have developed a variety of non-rotating Lorentz force flowmeters that can be used for a wide range of fluids and flow rates [8–11]. However, all of these flowmeters require calibration involving either analytical and/or numerical modeling [12–14], external calibration equipment [15, 16], or redundant flowmeters [4, 17]. This need for calibration introduces inconvenience and added expense that prevents Lorentz-force velocimetry from being more widely adopted in a range of industrial applications.

Parts of this section have been adapted from: M G Hvasta et al 2017 Meas. Sci. Technol. 28 085901.
To address these problems, this paper investigates the use of a new ‘weighted magnetic bearing’ (WMB) within an RLFF. It was found that the new bearing reduces the impact of friction on RLFF performance to negligible levels. From a practical standpoint, there were two major benefits to operating under ‘frictionless’ conditions. First, the RLFF was able to measure low flow rates that could not be detected when conventional bearings were used [5]. Second, since the trivially small frictional forces did not meaningfully affect the motion of the RLFF discs, the magnets moved at the average velocity of the flow, as predicted by Bucenieks [18]. This enables an RLFF to make accurate volumetric flow measurements without any calibration or prior knowledge of the fluid properties.

This paper is organized as follows. An overview of RLFF theory is presented in section 2, where attention is focused on the impact of the fluid velocity profile and the RLFF equations of motion. The different loss mechanisms inherent to the RLFF used in this experiment are described in section 3. It is emphasized that frictional losses in the bearing can be minimized while other loss mechanisms, such as windage losses, are only negligibly small at low angular velocities. Accordingly, the RLFF used in this experiment can only be expected to operate without calibration for low flow rates. The experimental setup used to test WMB operation is described in section 4. Experimental results comparing the calibrationless RLFF to a commercially available flowmeter can be found in section 5. The experiments showed that the RLFF flow measurements closely agreed with the commercially available flowmeter. A discussion of the results and a possible path forward for RLFF research is given in section 6.

2. RLFF theory

2.1. Impact of velocity profile

As explained in previous works [4, 15, 19], the Lorentz drag force ($F_D$) between a flowing, electrically-conductive liquid and the RLFF magnets can typically be described as

$$F_D \propto \sigma (v_0 - v_{mag}) B^2$$

where $\sigma$ is the electrical conductivity of the fluid, $v_0$ is the average velocity of the fluid, $v_{mag}$ is the average magnet velocity, $B$ is the magnitude of the externally applied magnetic field, and $v_{rel}$ is the relative velocity between the magnets and the fluid.

There are two notable situations where equation (1) may not yield accurate results. First, for equation (1) to be valid the magnitude of the induced or ‘secondary’ magnetic field ($b$) must be negligible compared to the external magnetic field ($b << B$) [20, 21]. This condition is satisfied if the magnetic Reynolds number ($\text{Re}_m$) is very small. For an RLFF, $\text{Re}_m$ can be calculated as

$$\text{Re}_m = \frac{\text{Magnetic Induction}}{\text{Magnetic Diffusion}} = \frac{v_{rel}D_h\mu_0 \sigma}{\nu}$$

where $D_h$ is the hydraulic diameter of the flow. It is conceivable that $\text{Re}_m$ could become quite large in fast-flowing systems where friction or other loss mechanisms prevent the RLFF magnets from moving at velocities comparable to $v_0$.

Figure 1. A depiction of a RLFF. In the ‘side view’, the flow is directed out of the page. Adapted from [5]. © 2017 Not subject to copyright in the USA. Contribution of U.S. Department of Energy.
However for this work and other similar research projects dealing with electrically conductive flows through pipes or tubes, numerous experiments have demonstrated the validity of equation (1) and verified the implicit assumption that Lorentz-force flowmeters are insensitive to different velocity profiles [5, 12, 18, 22]. In multiple studies, equation (1) has been experimentally confirmed by installing adjustable flow blockages just upstream of Lorentz-force flowmeters and then investigating the changes in the measured output. In these experiments it was found that the Lorentz-force flowmeter measurements were not affected by the changes in the velocity profile, especially at lower flow rates [4, 6, 11, 16]. Similarly, equation (1) has also been validated through experiments where flowing liquid systems were modeled or calibrated using uniform velocity profiles (slug-flow). The accuracy of the ‘slug-flow’ approximation was demonstrated numerically in turbulent salt-water flows [10] and experimentally in calibration facilities where solid metal proxies are used to calibrate Lorentz-force flowmeters [15]. In general, it is acknowledged that the indifference of Lorentz-force flowmeters towards the peculiarities of a given velocity profile make them useful and versatile for volumetric flow rate measurements [4, 18].

Nonetheless, it is recommended that flowmeter best-practices are followed whenever installing RLFFs (e.g. avoid bubbles/entrained gas, ensure the duct is completely filled, use flow conditioners if required, etc) [23–25]. If operators believe transient or otherwise atypical velocity profiles are a concern during flow measurement, a hydrodynamic entrance length (L) of approximately ten hydraulic diameters has been found to provide a nearly developed velocity profile for many applications [26]. Otherwise, the hydrodynamic entrance length for specific flows can be calculated using correlations below [26, 27]:

\[
\begin{align*}
\text{Laminar: } \frac{L}{D_h} &= 0.05 \text{Re} = 0.05 \left( \frac{\rho \omega v_0}{\mu} \right) \\
\text{Turbulent: } \frac{L}{D_h} &= 1.359 \text{Re}^{0.25} = 1.359 \left( \frac{\rho \omega v_0}{\mu} \right)^{0.25}
\end{align*}
\]

where \(\mu\) and \(\rho\) are the viscosity and density of the fluid respectively.

2.2. Flowmeter equations of motion

When an RLFF is used to measure the velocity of an electrically-conductive fluid moving through a duct that is both electrically-insulating and non-magnetic, the total torque on the flowmeter can be described as [5]

\[
\sum \tau = \tau_L [v_0 - \omega \cdot r] + \tau_F [\omega] = I \alpha
\]

\[
\sum \tau = K_L (v_0 - \omega \cdot r) + (\tau_B [\omega] + \tau_W [\omega]) = I \alpha
\]

where \(\tau_L\) is the torque generated by the Lorentz drag force \((F_B)\), \(\omega\) is the angular velocity of the RLFF, \(r\) is the effective radius of the RLFF, \(\tau_F\) is the net frictional torque resulting from the combined losses in the bearings \((\tau_B)\) and air resistance \((\tau_W)\), \(I\) is the moment of inertia of the RLFF, and \(\alpha\) is the angular acceleration of the RLFF. The constant \(K_L\) accounts for the fluid, magnetic, and geometric properties of a particular RLFF setup.

Under idealized conditions, where there are no frictional losses in the flowmeter \((\tau = 0)\), equation (4) shows that the magnet velocity would match the fluid velocity during steady-state operation \((\alpha = 0)\), namely

\[
v_0 = \omega r.
\]

In practice, a real flowmeter will always have some frictional losses, but the goal of this study is to design and build a calibrationless RLFF that minimizes these losses to such an extent that assuming \(\tau = 0\) leads to acceptably small errors in flow measurements. The next section will describe ways to quantify and minimize forces that oppose the rotation of an RLFF.

3. Bearing and windage losses

3.1. Bearing losses

For this experiment the losses in the bearing can be categorized as either frictional losses \((P_F)\) or eddy current losses \((P_{EC})\).

3.1.1. Frictional losses. Frictional losses \((P_F)\) in a rotating body can be described using the following equations [28–30]:

\[
P_F = \tau_l \omega
\]

\[
\tau_l = \int_0^r \mu_k [r] F_N [r] \, dr
\]

where \(\tau_l\) is the frictional torque, \(r\) is the lever-arm associated with the frictional forces, \(\mu_k\) is the coefficient of kinetic friction, and \(F_N\) is normal force between the sliding surfaces within a bearing.

The WMB used in this experiment is depicted in figure 2. The WMB reduces \(P_F\) in three ways. First, \(\mu_k\) is minimized by carefully selecting sliding surfaces to be hard and smooth. The coefficient of kinetic friction is further reduced by lubricating the sliding surfaces. Secondly, the lever-arm that the frictional forces use to act upon the rotating body is minimized by centering the connection between the two spheres directly above the RLFF. In theory the contact area between two spheres is a vanishingly small point, but in practice, because the magnet and ball bearing are imperfect spheres made of non-idealized materials, \(r \tau \gg 0\). Lastly, the normal force is minimized by offsetting the magnetic force \((F_B)\) with a counter-weight in such a way that the gravitational force \((F_g)\) is approximately equal to the magnetic force. More specifically, the magnitude of the normal force within a WMB, such as the one shown in figure 2, can be calculated using the following force-balance:

\[
F_N = F_B - F_g = F_B - mg
\]
where $m$ is the combined mass of the RLFF and the counter-weight, and $g \approx 9.81 \text{ (m s}^{-2})$. For a given magnetic force $F_B$, the mass of a counter-weight can be increased until the normal force is arbitrarily small ($F_N > 0$ for the bearing to stay connected). The magnitude of $F_B$ for any magnet configuration can be numerically calculated using the techniques described by Meeker [31].

If a WMB is used in a single-sided RLFF, such as those used by Priede et al [32] or Bucenieks [18, 22], care should be given to account for electromagnetic lift forces generated by the RLFF [33–35]. These lift forces ($F_L$) could be problematic during flowmeter startup or flow transients when the relative velocity between the fluid and the magnets is the greatest. Depending on the geometry of the single-sided RLFF, $F_L$ could either increase the normal force and friction between the magnet and the ball bearing or cause the magnet and ball bearing to separate ($F_N = 0$) in very carefully balanced WMBs. As described by Reitz and Davis [36, 37] the ratio of drag force ($F_D$—see equation (1)) to $F_L$ can be described as

$$\frac{F_D}{F_L} \sim \frac{2}{v_{rel} \sigma \mu_0 D_h^3}. \quad (9)$$

So, during steady-state operation when $v_{rel} \approx 0$, the lift forces will be minimized.

### 3.2. Eddy current losses

Due to the design of the WMB used in this experiment, electrically conductive components rotate within a magnetic field. Therefore, it is important to investigate the impact of eddy current losses ($P_{EC}$) within the bearing. As found in previous work, the eddy current losses can be described as [38–40]

$$P_{EC} \propto \sigma B^2 \omega^2. \quad (10)$$

At low flow rates where the corresponding angular velocity of the RLFF is small, the impact of $P_{EC}$ quickly diminishes. At higher flow rates, the impact of eddy current losses can be reduced by constructing the WMB from materials that have high electrical resistivities. Similarly, rotating components within the magnetic field could be fabricated using thin laminations of magnetic materials to help reduce $P_{EC}$ [41–43]. It is also possible that an electromagnet could be used within the WMB instead of a permanent magnet. In this case, the magnitude of the magnetic field could be adjusted to be as low as possible in order to minimize the effects of the eddy current losses.

### 3.3. Windage losses

Windage losses ($P_W$) are due to air resistance on the rotating parts of the RLFF. For simple cylindrical geometries like those found in most RLFFs, it has been shown that $P_W$ scales as [44, 45]

$$P_W \propto D^3 \omega^3. \quad (11)$$

Windage losses could have a substantial impact on RLFF performance at high rotational velocities. One way to reduce the impact of these losses is to design an RLFF to be more aerodynamic, which could be accomplished by eliminating bluff surfaces and keeping the overall size for the RLFF as small as practical. Alternatively, since windage losses become exceedingly small as the angular velocity approaches zero, RLFFs also can be designed to operate under low-flow conditions where the windage losses are negligible.
4. Experimental setup

The double-sided RLFF used in this experiment was constructed with a WMB, as pictured in figure 3. The RLFF magnet assembly consisted of two aluminum discs ($D = 25.4$ (cm)) each containing eight evenly spaced NdFeB N42 magnets ($5.08 \times 2.54 \times 1.27$ (cm)). The discs were arranged to produce an alternating magnetic field across the fluid. Thirty evenly-spaced optical markers along the rim of the top disc allowed the angular velocity of the disc to be measured with an optical tachometer. The output from the tachometer was collected using a LabVIEW-based data acquisition system. (This configuration closely resembles the RLFF setup used in previous work by this research group [5].)

An AISI 52100 chrome steel ball bearing with a diameter of 2.54 (cm) was attached to the top of the RLFF shaft. The RLFF assembly was then suspended from a spherical, nickel-plated NdFeB magnet with a diameter of 2.54 (cm). Prior to operation, a light coating of WD-40 was applied to the sliding surfaces of the ball bearing and magnet with a lint free cloth. Annular counter-weights were also concentrically fitted onto the central shaft of the RLFF. The counter-weights were designed to be easily accessible so that the total mass of the flowmeter could be adjusted by adding or removing discs. The entire device was supported by a frame that allowed the RLFF to be positioned around the liquid-metal tube. The frame was made from wood in order to avoid interactions with the RLFF magnets (e.g. magnetic attraction, induced eddy currents, etc).

The RLFF design was tested within the Liquid Metal eXperiment Upgrade (LMX-U) at Princeton Plasma Physics Laboratory (PPPL), which has been previously described in other works [46, 47]. LMX-U uses a gear-pump to circulate a liquid metal known as galinstan (Ga$_{67}$In$_{20.5}$Sn$_{12.5}$ wt.%) through plastic tubing with an inner diameter of $d = 3.97$ (cm). The electrical conductivity of galinstan is approximately $3.1 \times 10^6$ (S m$^{-1}$) [48–50]. A commercially available Omega FMG96 electromagnetic flowmeter (EMFM) was used to verify the results of the RLFF. The certified FMG96 was factory calibrated and tested according to NIST standards. This electromagnetic flowmeter has a rated accuracy of $\pm 5\%$ at flow rates of interest to this paper ($4.9–12.5$ (l min$^{-1}$)).

For this experiment the range of measured average velocities was $v_0 \approx 0.04–0.15$ (m s$^{-1}$) ($Re \approx 4.3E3–1.6E4$ [–]). Using equation (3), the turbulent hydrodynamic entrance length was calculated to be a maximum of 15 hydraulic diameters. During testing, the RLFF was installed approximately 18 hydraulic diameters downstream of the pump and all other flow obstructions, so it was assumed that the velocity profile was fully developed during all tests.

5. Results

5.1. Use of counterweights to reduce frictional losses

Frictional losses within the WMB were investigated for different counter-weight masses. As shown in figure 4, increasing the mass of the counter-weight reduced the normal force between the spherical magnet and the ball bearing thereby reducing frictional losses, as predicted by equations (7) and (8). The deceleration curves shown in figure 4 were measured for different counter-weight masses while the RLFF was positioned away from the liquid-metal or any other external conductive or magnetic materials ($\tau_L = 0$).

Without a counter-weight, the RLFF took about 20 (min) to decelerate from an initial angular velocity of approximately...
3.2 \text{ (rad s}^{-1})\), which was already a large improvement over previous work [5]. With a 1473 (g) counterweight, the RLFF spun freely for over 120 (min) when given the same initial angular velocity.

5.2. Flowmeter results

The frictionless assumption \((\tau_F = 0)\) was tested by comparing uncalibrated RLFF flow rate measurements to the output of the FMG96, as shown in figure 5. The uncalibrated RLFF values were calculated using equation (5), where the effective radius \((r \approx 9.84 \text{ (cm)})\) was assumed to be the distance from the RLFF axis of rotation to the center of the pipe. It was determined that the uncertainty in the RLFF measurements was about 3.3\%, which primarily resulted from errors in measuring the effective radius since \(\omega\) was very accurately known. This error corresponds to the effective radius being known to within approximately \(\pm 1.5\) (mm). The average flow velocity of the galinstan was also calculated using the volumetric flow rate \((Q)\) data measured by the EMFM and the following relation:

\[
\quad v_0 = \frac{4 Q}{\pi d^2},
\]

Figure 5 shows that the uncalibrated results from the RLFF fall within the \(\pm 5\%\) uncertainty margin of the EMFM. The RLFF was tested on two separate dates by two different operators. The device was removed and reinstalled around the liquid-metal filled pipe between each trial. Due to the RLFF not being installed in the exact same position during each test, small differences in the two data sets are expected.

It is also noteworthy that the RLFF was capable of generating steady outputs at flow rates lower than what the EMFM could measure. During the experiment the gear pump could operate smoothly at speeds as low as 50 (RPM). Assuming there was no fluid slipping past the gears in the pump, which experiments have shown to be an overly optimistic assumption, the maximum expected output of the pump would be \(\approx 2.50 \text{ (1 min}^{-1})\) or \(\approx 3.37 \text{ (cm}^{-1})\) [52]. Under these conditions the RLFF rotated at \(\omega \approx 0.33 \text{ (rad s}^{-1})\) which corresponded to a flow rate of \(\approx 3.06 \text{ (cm}^{-1})\). This is a promising and encouraging result because it is close to but less than the maximum expected output of the gear pump. Based on these results, operational experience with the RLFF, and the deceleration curves in figure 4, it is extremely likely that the RLFF would be able to measure even smaller flow rates if it were tested with different experimental hardware.

6. Discussion and future work

A novel RLFF was designed, built, and tested. The performance of the WMB was adjusted by changing the mass of a counter-weight in order to reduce the normal force between the spherical magnet and the steel ball bearing. For carefully selected counter-weight masses, it was shown that a WMB can nearly eliminate frictional losses in an RLFF.

The advantages of using a WMB to minimize frictional losses in an RLFF are two-fold:

- The RLFF can be used without calibration and still achieve accuracies comparable to commercially available flowmeters.
- The RLFF has improved sensitivity at low flow rates when compared to using traditional roller bearings [5]. (For this experiment, the RLFF could measure flow rates that were too small for the commercially available flowmeter used to verify RLFF operation.)

These improvements have the potential to make RLFFs better suited to industrial applications where easy, precise, and accurate measurements of flows are required. However, it should be emphasized that not all loss mechanisms within the RLFF were reduced (see section 3). At large RLFF angular velocities, windage or eddy current losses could become appreciable and require the device to be calibrated. Furthermore, it should be reiterated that the fluid in this experiment flowed through non-magnetic, electrically-insulating plastic tubing. The use of metal pipes or tubes would have likely affected the output of the RLFF [3, 18, 22].

Work on RLFFs will continue at Princeton University with a focus on the following areas:

- The development of lighter components to improve the transient performance and response time of the flowmeter.
- Better techniques to install and align the RLFF with the liquid-metal filled tube. More accurate and consistent positioning of the flowmeter around the pipe would reduce uncertainty in the measurements.
- Streamlining the RLFF magnet assembly in order to reduce windage losses and allow for more accurate measurements at higher flow rates.

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The digital data for this paper can be found at http://arks.princeton.edu/ark:/88435/dsp01x920g025r
Appendix

A.1. Nomenclature

<table>
<thead>
<tr>
<th>Variable</th>
<th>Meaning</th>
<th>Units</th>
</tr>
</thead>
<tbody>
<tr>
<td>b</td>
<td>Induced magnetic field</td>
<td>T</td>
</tr>
<tr>
<td>B</td>
<td>External magnetic field</td>
<td>T</td>
</tr>
<tr>
<td>d</td>
<td>Inner diameter of pipe or tube</td>
<td>m</td>
</tr>
<tr>
<td>D</td>
<td>Diameter of RLFF disc</td>
<td>m</td>
</tr>
<tr>
<td>D_h</td>
<td>Hydraulic diameter</td>
<td>m</td>
</tr>
<tr>
<td>F</td>
<td>Force</td>
<td>N</td>
</tr>
<tr>
<td>g</td>
<td>Gravitational acceleration</td>
<td>m s⁻²</td>
</tr>
<tr>
<td>I</td>
<td>Moment of inertia</td>
<td>kg m⁻²</td>
</tr>
<tr>
<td>K_L</td>
<td>Lorentz factor</td>
<td>N s⁻¹</td>
</tr>
<tr>
<td>L</td>
<td>Hydrodynamic entrance length</td>
<td>m</td>
</tr>
<tr>
<td>m</td>
<td>Mass</td>
<td>kg</td>
</tr>
<tr>
<td>N</td>
<td>Interaction parameter</td>
<td>—</td>
</tr>
<tr>
<td>P</td>
<td>Power loss</td>
<td>W</td>
</tr>
<tr>
<td>Q</td>
<td>Volumetric flow rate</td>
<td>m³ s⁻¹</td>
</tr>
<tr>
<td>r</td>
<td>Radius</td>
<td>m</td>
</tr>
<tr>
<td>Re</td>
<td>Reynolds #</td>
<td>—</td>
</tr>
<tr>
<td>Re_m</td>
<td>Magnetic Reynolds #</td>
<td>—</td>
</tr>
<tr>
<td>v₀</td>
<td>Average fluid velocity</td>
<td>m s⁻¹</td>
</tr>
<tr>
<td>v avg</td>
<td>Average magnet velocity</td>
<td>m s⁻¹</td>
</tr>
<tr>
<td>v rel</td>
<td>Relative velocity</td>
<td>m s⁻¹</td>
</tr>
<tr>
<td>α</td>
<td>Angular acceleration</td>
<td>rad s⁻²</td>
</tr>
<tr>
<td>µ</td>
<td>Viscosity</td>
<td>Pa s⁻¹</td>
</tr>
<tr>
<td>µ_0</td>
<td>Vacuum permeability</td>
<td>H m⁻¹</td>
</tr>
<tr>
<td>µ_κ</td>
<td>Coefficient of kinetic friction</td>
<td>—</td>
</tr>
<tr>
<td>ρ</td>
<td>Density</td>
<td>kg/m³</td>
</tr>
<tr>
<td>σ</td>
<td>Electrical conductivity</td>
<td>S m⁻¹</td>
</tr>
<tr>
<td>τ</td>
<td>Torque</td>
<td>N m</td>
</tr>
<tr>
<td>ω</td>
<td>Angular velocity</td>
<td>rad s⁻¹</td>
</tr>
</tbody>
</table>

Abbreviation | Meaning
---|---
EMFM | Electromagnetic flowmeter
LMX-U | Liquid Metal eXperiment—Upgrade
PPPL | Princeton Plasma Physics Laboratory
RLFF | Rotating Lorentz-force flowmeter
WMB | Weighted magnetic bearing

Mathematical convention | Meaning
---|---
A(x) | ‘A’ is a function of ‘x’

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